



The global attractors and their Hausdorff and fractal dimensions estimation for the higher-order nonlinear Kirchhoff-type equation with nonlinear strongly damped terms*

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Abstract

In this paper, we study the long time behavior of solution to the initial boundary value problems for higher -order kirchhoff-type equation with nonlinear strongly dissipation:

$$u_{tt} + (-\Delta)^m u_t + (\int_{\Omega} |\nabla^m u|^2)^q (-\Delta)^m u + h(u_t) = f(x).$$

At first, we prove the existence and uniqueness of the solution by priori estimate and Galerkin method then we establish the existence of global attractors, at last, we consider that estimation of upper bounds of Hausdorff and fractal dimensions for the global attractors are obtain.

Keywords: Higher-order nonlinear Kirchhoff wave equation; The existence and uniqueness; The Global attractors; Hausdorff dimensions; Fractal dimensions

1 Introduction

In this paper we concerned with the long time behavior of solution to the initial boundary value problems for Higher-order Kirchhoff-type equation with nonlinear strongly dissipation :

$$u_{tt} + (-\Delta)^m u_t + (\int_{\Omega} |\nabla^m u|^2)^q (-\Delta)^m u + h(u_t) = f(x) \quad (1.1)$$

$$u(x, t) = 0, \frac{\partial^i u}{\partial v^i} = 0, \quad i = 1, 2, \dots, m-1, x \in \partial\Omega, t > 0, \quad (1.2)$$

$$u(x, 0) = u_0, u_t = u_1(x), \quad x \in \partial\Omega \quad (1.3)$$

Where $\Omega \subset \mathbb{R}^2$ is bounded open domain with smooth boundary; v is the outer norm vector; $m > 1$ is a positive integer, and $q > 0$ is a positive constants, $h(u_t)$ is a nonlinear forcing, $(-\Delta)^m u_t$ is a strongly dissipation.

There have been many researches on the well-positive and the longtime dynamics for Kirchhoff equation. we can see [1-6], FUCAI Li [5] deals with the higher-order kirchhoff-type equation with nonlinear dissipation:

$$u_{tt} + (\int_{\Omega} |\nabla^m u|^2)^q (-\Delta)^m u + u_t |u_t|^r = |u|^p u, \quad x \in \Omega, \quad t > 0, \quad (1.4)$$



$$u(x,t) = 0, \frac{\partial^i u}{\partial v^i} = 0, \quad i = 1, 2, \dots, m-1, x \in \partial\Omega, t > 0, \quad (1.5)$$

$$u(x,0) = u_0, u_t = u_1(x), \quad x \in \partial\Omega. \quad (1.6)$$

In a bounded domain, where $m > 1$ is a positive integer, $p, q, r > 0$ are positive constants and obtain that the solution exists globally if $p \leq r$, while if $p > \max\{r, 2q\}$, then for any initial data with negative initial energy, the solution blows up at finite time in L^{p+2} norm.

Yang Zhijian, Wang Yunqing [6] also studied the global attractor for the Kirchhoff type equation with a strong dissipation:

$$u_{tt} - M(\|\Delta u\|^2) \Delta u - \Delta u_t + h(u_t) + g(u) = f(x) \quad \text{in } \Omega \times \mathbb{R}^+, \quad (1.7)$$

$$u(x,t)|_{\partial\Omega} = 0, \quad t > 0, \quad (1.8)$$

$$u(x,0) = u_0(x), \quad u_t(x,0) = u_1(x), \quad x \in \Omega. \quad (1.9)$$

Where $M(s) = 1 + s^{\frac{m}{2}}$, $1 \leq m \leq \frac{4}{N-2}$, Ω is a bounded domain in \mathbb{R}^N , with smooth boundary $\partial\Omega$, $h(s)$ and $g(s)$ are nonlinear functions, and $f(x)$ is an external force term. It proves that the relative continuous semigroup $S(t)$ possesses in the phase space with low regularity a global attractor which is connected.

Yang zhijian, Cheng Jianling [7] studies the asymptotic behavior of solutions to the Kirchhoff-type equation:

$$u_{tt} - M(\|\Delta u\|^2) \Delta u - \Delta u_t + h(u_t) + g(u) = f(x) \quad \text{in } \Omega \times \mathbb{R}^+, \quad (1.10)$$

$$u(x,t)|_{\partial\Omega} = 0, \quad t > 0, \quad (1.11)$$

$$u(x,0) = u_0(x), \quad u_t(x,0) = u_1(x), \quad x \in \Omega. \quad (1.12)$$

They prove that the related continuous semigroup $S(t)$ possesses in phase $X = (H^2(\Omega) \cap H_0^1)$

$\times H_0^1(\Omega)$ a global attractor. At the end of the paper, an example is shown, which indicates the

Existence of nonlinear functions $g(x, u)$ and $h(u_t)$.

Recently, Meixia Wang, Cuicui Tian, Guoguang Lin [8] studied the global attractor and dimension for a 2D generalized Anisotropy Kuramoto-Sivashinsky equation:



$$u_t + \alpha \Delta^2 u + \gamma u + (\varphi(u))_{xx} + (g(u))_{yy} = f(x), \quad (x, y) \in \Omega \subset R^2, \quad (1.13)$$

$$u(x, y, t)|_{t=0} = u_0(x), \quad (x, y) \in \Omega \subset R^2, \quad (1.14)$$

$$u(x, y, t)|_{\partial\Omega=0} = 0, \quad \Delta u(x, y, t)|_{\partial\Omega} = 0, \quad (x, y) \in \Omega \subset R^2, \quad (1.15)$$

Where $\Omega \subset R^2$ is bounded set; $\partial\Omega$ is the bound of Ω ; $\varphi(u)$ and $g(u)$ are considered as smooth function of $u(x, y, t)$. under the existence of the global solution, it proves that the global attractor and Hausdorff dimensions and fractal dimension.

The paper is arranged as follows.in section 2,we state some preliminaries under the assume of Lemma 1and Lemma 2, we get the existence and uniqueness of solution; in section 3,we obtain the global attractors for the problems (1.1)-(1.3);in section 4.we consider that the global attractor of the the above mentioned problems (1.1)-(1.3) has finite Hausdorff dimensions and fractal dimensions.

2 Preliminaries

For convenience,we denote the norm and scalar product in $L^2(\Omega)$ by $\|\cdot\|$ and (\cdot, \cdot) ; $f = f(x)$,

$$H^k = H^k(\Omega), \quad H_0^k = H_0^k(\Omega), \quad H^{-k} = H^{-k}(\Omega), \quad \|\cdot\| = \|\cdot\|_{L^2}, \quad C_i (i = 0 \dots 8), \quad \kappa \quad \text{are constants},$$

$$K_0 \geq \max\left\{\frac{q\varepsilon}{q+1}, \frac{q}{q+1}, \frac{2\alpha\varepsilon^2}{q+1}, \frac{4C_1\varepsilon(2q-2p+2)}{2p+2}\right\}.$$

In this section, we present some materials needed in the proof of our results, state a global existence result, and prove our main result. For this reason, we assume that and notations needed in the proof of our results.For this reason, we assume that

(G_1) there exist

$$0 < \delta < \frac{1}{2}, \quad \|h(u_t)\|_{H^{-m}} \leq C_0(h(u_t, u))^{1-\delta}, \quad \forall u \in H_0^m,$$

(G_2) there exist constant

$$0 < \sigma < 1, \quad \|h(v)\| \leq C_1(R)(1 + \|\Delta v\|)^{1-\sigma}, \quad \forall v \in H^{2m} \cap H_0^m, \quad \|v\| \leq R,$$

(G_3) $|h'(s)| \leq C_2$

Lemma 1. Assume (G_1) hold, and $(u_0, u_1) \in H^m(\Omega) \times L^2(\Omega)$, $f \in L^2(\Omega)$, $v = u_t + \varepsilon u$,then the solution



$(u, v) \in H^m(\Omega) \times L^2(\Omega)$, and

$$\begin{aligned} \| (u, v) \|_{H^m \times L^2}^2 &= \| \nabla^m u \|_{}^2 + \| v \|_{}^2 \leq \frac{W(0) e^{-\alpha_0 t}}{1 - \varepsilon} + \frac{\frac{C_3 (1 - e^{-\alpha_0 t})}{\alpha_0} - K_0 + \frac{q}{q+1}}{1 - \varepsilon} \quad (2.1) \quad \text{Where } v = u_t + \varepsilon u, \\ 0 < \varepsilon < \min\{1, \frac{\sqrt{1 + 4\lambda_1^m} - 1}{2}, \frac{q+1}{4C_1 p}, \frac{q+1}{2\alpha}\} \quad W(0) &= \| v_0 \|_{}^2 + \end{aligned}$$

$$\frac{q}{q+1} \| \nabla^m u \|_{}^{2q+2} - \varepsilon \| \nabla^m u \|_{}^2 + K_0, \quad v_0 = u_1 + \varepsilon u_0, \text{ thus there exist } R_0 \text{ and } t_1 = t_1(\Omega) > 0,$$

Such that

$$\| (u, v) \|_{H^m \times L^2} = \| \nabla^m u \|_{}^2 + \| v \|_{}^2 \leq R_0 (t > t_1).$$

Proof. Let $v = u_t + \varepsilon u$ we multiply v with both sides of equation (1.1) and obtain

$$(u_u + (-\Delta)^m u_t + (\int_{\Omega} |\nabla^m u|^2)^q (-\Delta)^m u + h(u_t), v) = (f(x), v), \quad (2.2)$$

$$(u_u, v) = (v_t - \varepsilon u_t, v)$$

$$\begin{aligned} &= (v_t, v) - \varepsilon (v - \varepsilon u, v) \\ &= \frac{1}{2} \frac{d}{dt} \| v \|_{}^2 - \varepsilon (v - \varepsilon u, v) \quad (2.4) \end{aligned}$$

$$\geq \frac{1}{2} \frac{d}{dt} \| v \|_{}^2 - \varepsilon \| v \|_{}^2 - \frac{\varepsilon^2}{2\lambda_1^m} \| \nabla^m u \|_{}^2 - \frac{\varepsilon^2}{2} \| v \|_{}^2.$$

$$((- \Delta)^m u_t, v) = ((-\Delta)^m (v - \varepsilon u), v) \quad (2.5)$$

$$\geq \| \nabla^m v \|_{}^2 - \frac{1}{2} \frac{d}{dt} \| \nabla^m u \|_{}^2 - \varepsilon^2 \| \nabla^m u \|_{}^2$$

$$(\| \nabla^m u \|_{}^{2q} (-\Delta)^m u, v) = (\| \nabla^m u \|_{}^{2q} (-\Delta)^m u, u_t + \varepsilon u) \quad (2.6)$$

$$= \frac{1}{2q+1} \frac{d}{dt} \| \nabla^m u \|_{}^{2q+2} + \varepsilon \| \nabla^m u \|_{}^{2q+2}.$$

$$(h(u_t), v) = (h(u_t), u_t + \varepsilon u) \quad (2.7)$$



$$= (h(u_t), u_t) + (h(u_t), \varepsilon u).$$

$$\begin{aligned} \varepsilon |(h(u_t), u)| &\leq \|h(u_t)\|_{H^{-m}} \|\nabla^m u\| \\ &\leq \varepsilon C_0 (h(u_t), u_t)^{1-\delta} \|\nabla^m u\| \\ &\leq \frac{1}{2} (h(u_t), u_t) + C_1 \varepsilon^{2p} \|\nabla^m u\|^{2p} \quad (2.8) \\ &\leq \frac{1}{2} (h(u_t), u_t) + C_1 \varepsilon^2 \|\nabla^m u\|^{2p}. \end{aligned}$$

$$(f(x), v) \leq \|f\| \|v\| \leq \frac{\varepsilon^2}{2} \|v\|^2 + \frac{1}{2\varepsilon^2} \|f\|^2 \quad (2.9)$$

For above ,we have

$$\begin{aligned} \frac{1}{2} \frac{d}{dt} (\|v\|^2 + \frac{1}{q+1} \|\nabla^m u\|^{2q+2} - \varepsilon \|\nabla^m u\|^2) + \|\nabla^m u\|^2 + \varepsilon \|v\|^2 - \varepsilon^2 \|v\|^2 \\ + \varepsilon \|\nabla^m u\|^{2q+2} - (\varepsilon^2 + \frac{\varepsilon^2}{2\lambda_1}) \|\nabla^m u\|^2 - C_1 \varepsilon^2 \|\nabla^m u\|^{2p} \quad (2.10) \\ \leq \frac{1}{2\varepsilon^2} \|f\|^2. \end{aligned}$$

By using Poincare inequality,we obtain: $\|\nabla^m v\|^2 \geq \lambda_1^m \|v\|^2$, then we have

$$\begin{aligned} \frac{d}{dt} (\|v\|^2 + \frac{1}{q+1} \|\nabla^m u\|^{2q+2} - \varepsilon \|\nabla^m u\|^2 + K_0) + (2\lambda_1 + 2\varepsilon - 2\varepsilon^2) \|v\|^2 \\ + 2\varepsilon \|\nabla^m u\|^{2q+2} - (2\varepsilon^2 + \frac{\varepsilon^2}{\lambda_1}) \|\nabla^m u\|^2 - 2C_1 \varepsilon^2 \|\nabla^m u\|^{2p} + 2K_0 \quad (2.11) \\ \leq \frac{1}{\varepsilon^2} \|f\|^2. \end{aligned}$$

By using Young's inequality,we obtain

$$\|\nabla^m u\|^2 \leq \frac{1}{q+1} \|\nabla^m u\|^{2q+2} + \frac{q}{q+1}, \quad (2.12)$$

$$- \varepsilon \|\nabla^m u\|^2 \geq - \frac{\varepsilon}{q+1} \|\nabla^m u\|^{2q+2} - \frac{q\varepsilon}{q+1}, \quad (2.13)$$

Then we have



$$\|v\|^2 + \frac{1}{q+1} \|\nabla^m u\|^{2q+2} - \frac{\varepsilon}{q+1} \|\nabla^m u\|^{2q+2} + K_0 \geq 0. \quad (2.14)$$

Taking $\alpha = 2 + \frac{1}{\lambda_1^m}$, and $p \leq q+1$ using Young's inequality we obtain

$$-\alpha\varepsilon^2 \|\nabla^m u\|^2 \geq -\frac{\alpha\varepsilon^2}{q+1} \|\nabla^m u\|^{2q+2} - \frac{\alpha\varepsilon^2}{q+1}, \quad (2.15)$$

$$\|\nabla^m u\|^{2p} \leq \frac{2p}{2q+2} \|\nabla^m u\|^{2q+2} + \frac{2q-2p+2}{2p+2}, \quad (2.16)$$

$$-2C_0\varepsilon^2 \|\nabla^m u\|^{2p} \leq \frac{4pC_0\varepsilon^2}{2q+2} \|\nabla^m u\|^{2q+2} + \frac{2C_0\varepsilon^2(2q-2p+2)}{2p+2}. \quad (2.17)$$

So, we obtain

$$\frac{\varepsilon}{2} \|\nabla^m u\|^{2q+2} - \alpha\varepsilon^2 \|\nabla^m u\|^2 + \frac{K_0}{2} \geq \left(\frac{\varepsilon}{2} - \frac{\alpha\varepsilon^2}{q+1}\right) \|\nabla^m u\|^{2q+2} + \frac{K_0}{2} - \frac{\alpha\varepsilon^2}{q+1} \geq 0. \quad (2.18)$$

$$\begin{aligned} & \frac{\varepsilon}{2} \|\nabla^m u\|^{2q+2} - 2C_1\varepsilon^2 \|\nabla^m u\|^{2p} + \frac{K_0}{2} \\ & \geq \frac{\varepsilon}{2} \|\nabla^m u\|^{2q+2} - \frac{4C_1p\varepsilon^2}{2q+2} \|\nabla^m u\|^{2q+2} + \frac{K_0}{2} - \frac{2C_1\varepsilon^2(2q-2p+2)}{2p+2} \geq 0. \end{aligned} \quad (2.19)$$

So, we have

$$\begin{aligned} & \frac{d}{dt} (\|v\|^2 + \frac{1}{q+1} \|\nabla^m u\|^{2q+2} - \varepsilon \|\nabla^m u\|^2 + K_0) + (2\lambda_1 + 2\varepsilon - 2\varepsilon^2) \|v\|^2 \\ & + \varepsilon \|\nabla^m u\|^{2q+2} + K_0 \quad (2.20) \\ & \leq \frac{1}{\varepsilon^2} \|f\|^2 + 2K_0. \end{aligned}$$

Assume Next we take $\alpha_0 = \{2\lambda_1 - 2\varepsilon - 2\varepsilon^2, (q+1)\varepsilon, 1\}$, so we can obtain

$$\begin{aligned} & \frac{d}{dt} (\|v\|^2 + \frac{1}{q+1} \|\nabla^m u\|^{2q+2} - \varepsilon \|\nabla^m u\|^2 + K_0) \\ & + \alpha_0 (\|v\|^2 + \frac{1}{q+1} \|\nabla^m u\|^{2q+2} - \varepsilon \|\nabla^m u\|^2 + K_0) \quad (2.21) \end{aligned}$$



$$\leq \frac{1}{\varepsilon^2} \|f\|^2 + 2K_0.$$

Then we have

$$\frac{d}{dt}W(t) + \alpha_0 W(t) \leq C_3, \quad (2.22)$$

where $W(t) = \|v\|^2 + \frac{1}{q+1} \|\nabla^m u\|^{2q+2} - \varepsilon \|\nabla^m u\|^2 + K_0$, $C_3 = \frac{1}{\varepsilon^2} \|f\|^2 + 2K_0$, by using Gronwall inequality, we

obtain

$$W(t) \leq W(0)e^{-\alpha_0 t} + \frac{C_3(1 - e^{-\alpha_0 t})}{\alpha_0}, \quad (2.23)$$

From (2.12), we know

$$\|\nabla^m u\|^2 - \frac{q}{q+1} \leq \frac{1}{q+1} \|\nabla^m u\|^{2q+2}. \quad (2.24)$$

So

$$\|v\|^2 + \frac{1}{q+1} \|\nabla^m u\|^{2q+2} - \varepsilon \|\nabla^m u\|^2 + K_0 \geq \|v\|^2 + (1 - \varepsilon) \|\nabla^m u\|^2 + K_0 - \frac{q}{q+1} \geq 0, \quad (2.25)$$

From (2.23), we obtain

$$\|v\|^2 + (1 - \varepsilon) \|\nabla^m u\|^2 + K_0 - \frac{q}{q+1} \leq W(0)e^{-\alpha_0 t} + \frac{C_3(1 - e^{-\alpha_0 t})}{\alpha_0}, \quad (2.26)$$

Then we have

$$(1 - \varepsilon)(\|v\|^2 + \|\nabla^m u\|^2) \leq W(0)e^{-\alpha_0 t} + \frac{C_3(1 - e^{-\alpha_0 t})}{\alpha_0} - K_0 + \frac{q}{q+1}, \quad (2.27)$$

so, we obtain

$$\|(u, v)\|_{H^m \times L^2}^2 = \|\nabla^m u\|^2 + \|v\|^2 \leq \frac{W(0)e^{-\alpha_0 t}}{1 - \varepsilon} + \frac{\frac{C_3(1 - e^{-\alpha_0 t})}{\alpha_0} - K_0 + \frac{q}{q+1}}{1 - \varepsilon}. \quad (2.28)$$

Then

$$\overline{\lim}_{t \rightarrow \infty} \|(u, v)\|_{H^m \times L^2}^2 \leq \frac{\frac{C_3}{\alpha_0} - K_0 + \frac{q}{q+1}}{1 - \varepsilon}. \quad (2.29)$$

So, there exist R_0 and $t_0(\Omega) > 0$, such that



$$\| (u, v) \|_{H^m \times H^m}^2 = \| \nabla^m u \|_{}^2 + \| v \|_{}^2 \leq R_0 (t > t_1) \quad (2.30)$$

Lemma 2. In addition to the assumptions (G_2) hold, $f \in H^m(\Omega)$, then the solution (u, v) of the problems

(1.1)-(1.3) satisfies $(u, v) \in H^{2m}(\Omega) \times H^m(\Omega)$, and

$$\| (u, v) \|_{H^{2m} \times H^m}^2 = \| (-\Delta)^m u \|_{}^2 + \| \nabla^m v \|_{}^2 \leq \frac{Y(0)e^{-\beta_0 t}}{\beta_1} + \frac{C_4(1 - e^{-\beta_0 t})}{\beta_0 \beta_1} \quad (2.31)$$

In where $v = u_t + \varepsilon u$, and $Y(0) = (\| (-\nabla)^m u_0 \|^{2q} - \varepsilon) \| (-\Delta)^m u_0 \|_{}^2 + \| \nabla^m v_0 \|_{}^2$ thus there exist R_1 and $t_2 = t_2(\Omega) > 0$, such that

$$\| (u, v) \|_{H^{2m} \times L^2}^2 = \| (-\Delta)^m u \|_{}^2 + \| \nabla^m v \|_{}^2 \leq R_1 (t > t_1) \quad (2.32)$$

Proof. Let $(-\Delta)^m v = (-\Delta)^m u_t + \varepsilon (-\Delta)^m u$, we multiply $(-\Delta)^m v$ with both sides of equation (1.1), and obtain

$$(u_t + (-\Delta)^m u_t + (\int_{\Omega} |\nabla^m u|^2)^q (-\Delta)^m u + h(u_t), (-\Delta)^m v) = (f(x), (-\Delta)^m v), \quad (2.33)$$

$$(u_t, (-\Delta)^m v) \geq \frac{1}{2} \frac{d}{dt} \| \nabla^m v \|_{}^2 - \varepsilon \| \nabla^m v \|_{}^2 - \frac{\varepsilon^2}{2 \lambda_1} \| (-\Delta)^m u \|_{}^2 - \frac{\varepsilon^2}{2} \| \nabla^m v \|_{}^2 \quad (2.34)$$

$$((- \Delta)^m u_t, (-\Delta)^m v) = ((-\Delta)^m (v - \varepsilon u), (-\Delta)^m v) \quad (2.35)$$

$$\begin{aligned} &= \| (-\Delta)^m v \|_{}^2 - \frac{\varepsilon}{2} \frac{d}{dt} \| (-\Delta)^m u \|_{}^2 - \varepsilon^2 \| (-\Delta)^m u \|_{}^2 \\ &= \| \nabla^m u \|^{2q} (-\Delta)^m u, (-\Delta)^m v = \frac{1}{2} \frac{d}{dt} (\| \nabla^m u \|^{2q} \| (-\Delta)^m u \|_{}^2) - \frac{\| (-\Delta)^m u \|_{}^2}{2} \frac{d}{dt} \| \nabla^m u \|^{2q} \quad (2.36) \\ &\quad + \varepsilon \| \nabla^m u \|^{2q} \| (-\Delta)^m u \|_{}^2 \\ &\left| (h(u_t), (-\Delta)^m v) \right| \leq \frac{\| h(u_t) \|_{}^2}{2} + \frac{\| (-\Delta)^m v \|_{}^2}{2} \quad (2.37) \end{aligned}$$

From (G_2) , we have

$$\| h(u_t) \|_{}^2 \leq C_1^2 (R) (1 + \| (-\Delta)^m u_t \|_{}^{2(1-\sigma)}), \quad (2.38)$$

By using Young's inequality



$$\|h(u_t)\|^2 \leq \frac{\sigma}{\mu^{\frac{1}{\sigma}}} (C_1^2(R))^{\frac{1}{\sigma}} + (1-\sigma) \mu^{\frac{1}{1-\sigma}} ((1 + \|(-\Delta)^m u_t\|^{2(1-\sigma)})^{\frac{1}{1-\sigma}}), \quad (2.39)$$

and

$$\|h(u_t)\|^2 \leq C_5 + \frac{1}{4} \|(-\Delta)^m v\|^2 + \frac{\varepsilon^2}{4} \|(-\Delta)^m u\|^2, \quad (2.40)$$

Where $C_5 := \frac{\sigma}{\mu^{\frac{1}{\sigma}}} (C_1^2(R))^{\frac{1}{\sigma}} + 2(1-\sigma) \mu^{\frac{1}{1-\sigma}}$, we take proper μ , such that: $4(1-\sigma) \mu^{\frac{1}{1-\sigma}} = \frac{1}{4}$,

$$K_1 = C_5.$$

$$|(f(x), (-\Delta)^m v)| \leq \frac{\|\nabla^m f\|^2}{2\varepsilon^2} + \frac{\varepsilon^2 \|\nabla^m v\|^2}{2} \quad (2.41)$$

From above, we have

$$\begin{aligned} & \frac{d}{dt} (\|\nabla^m v\|^2 + \|\nabla^m u\|^{2q} \|(-\Delta)^m u\|^2 - \varepsilon \|(-\Delta)^m u\|^2) + \left(\frac{3\lambda_1^m}{4} - 2\varepsilon - 2\varepsilon^2 \right) \|\nabla^m v\|^2 \\ & + \left(-\frac{d}{dt} \|\nabla^m u\|^{2q} + 2\varepsilon \|\nabla^m u\|^{2q} - \frac{9\varepsilon^2}{4} - \frac{\varepsilon^2}{\lambda_1^m} \right) \|(-\Delta)^m u\|^2 \\ & \leq \frac{1}{\varepsilon^2} \|\nabla^m f\|^2 + 2K_1. \end{aligned} \quad (2.42)$$

Then we take proper ε , let $\frac{3\lambda_1^m}{4} - 2\varepsilon - 2\varepsilon^2 \geq 0$ and $\|\nabla^m u\|^{2q} - \varepsilon > 0$, next we assume exist a positive constant

$K > 0$, let $K - 2\varepsilon \geq 0$, satisfies

$$0 \leq K (\|\nabla^m u\|^{2q} - \varepsilon) \leq -\frac{d}{dt} \|\nabla^m u\|^{2q} + 2\varepsilon \|\nabla^m u\|^{2q} - \frac{9\varepsilon^2}{4} - \frac{\varepsilon^2}{\lambda_1^m} \quad (2.43)$$

where $C_6 := -\frac{9\varepsilon^2}{4} - \frac{\varepsilon^2}{\lambda_1^m} + K\varepsilon$ such that

$$(K - 2\varepsilon) \|\nabla^m u\|^{2q} + \frac{d}{dt} \|\nabla^m u\|^{2q} \leq C_6, \quad (2.44)$$

Multiplying (2.44) by $e^{(K-2\varepsilon)t}$ then

$$\|\nabla^m u\|^{2q} \frac{d}{dt} (e^{(K-2\varepsilon)t}) + e^{(K-2\varepsilon)t} \frac{d}{dt} \|\nabla^m u\|^{2q} \leq C_6 e^{(K-2\varepsilon)t}, \quad (2.45)$$

we integrate (2.45) with respect to time t and get that



$$\varepsilon < \left\| \nabla^m u \right\|^{2q} \leq \frac{C_6}{K - 2\varepsilon} (1 + \kappa e^{-(K-2\varepsilon)t}), \quad (2.46)$$

So,(2.43) exists a positive constant K .

Form above , we have

$$\begin{aligned} \frac{d}{dt} (\left\| \nabla^m v \right\|^2 + \left\| \nabla^m u \right\|^{2q} \left\| (-\Delta)^m u \right\|^2 - \varepsilon \left\| (-\Delta)^m u \right\|^2) + (\frac{3\lambda_1^m}{4} - 2\varepsilon - 2\varepsilon^2) \left\| \nabla^m v \right\|^2 \\ + K (\left\| \nabla^m u \right\|^{2q} - \varepsilon) \left\| (-\Delta)^m u \right\|^2 \\ \leq \frac{1}{\varepsilon^2} \left\| \nabla^m f \right\|^2 + 2K_1. \end{aligned} \quad (2.47)$$

he Taking $\beta_0 = \min\{ \frac{3\lambda_1^m}{4} - 2\varepsilon - 2\varepsilon^2, K \}$, $C_4 = \frac{1}{\varepsilon^2} \left\| \nabla^m f \right\|^2 + 2K_1$ then

$$\frac{d}{dt} Y(t) + \beta_0 Y(t) \leq C_4, \quad (2.48)$$

assumptions where $Y(t) = \left\| \nabla^m v \right\|^2 + (\left\| \nabla^m u \right\|^{2q} - \varepsilon) \left\| (-\Delta)^m u \right\|^2$ by using Gronwall inequality,then

$$Y(t) \leq Y(0) e^{-\beta_0 t} + \frac{C_4}{\beta_0} (1 - e^{-\beta_0 t}). \quad (2.49)$$

Let $\beta_1 = \min\{ 1, \inf_{t \geq 0} \left\| \nabla^m u \right\|^{2q} - \varepsilon \}$,we get $\beta_1 (\left\| \nabla^m v \right\|^2 + \left\| (-\Delta)^m u \right\|^2) \leq Y(t)$.

so we get

$$\left\| (u, v) \right\|_{H^{2m} \times H^m}^2 = \left\| (-\Delta)^m u \right\|^2 + \left\| \nabla^m v \right\|^2 \leq \frac{Y(0) e^{-\beta_0 t}}{\beta_1} + \frac{C_4 (1 - e^{-\beta_0 t})}{\beta_0 \beta_1}. \quad (2.50)$$

where $Y(0) = (\left\| (-\Delta)^m u_0 \right\|^{2q} - \varepsilon) \left\| (-\Delta)^m u_0 \right\|^2 + \left\| \nabla^m v_0 \right\|^2$, $u_0 = u_1 + \varepsilon u_0$ then

$$\overline{\lim}_{t \rightarrow \infty} \left\| (u, v) \right\|_{H^{2m} \times H^m}^2 \leq \frac{C_4}{\beta_0 \beta_1}. \quad (2.51)$$

So,there exist R_1 and $t_1 = t_1(\Omega) > 0$, such that

$$\left\| (u, v) \right\|_{H^{2m} \times H^m}^2 = \left\| (-\Delta)^m u \right\|^2 + \left\| \nabla^m v \right\|^2 \leq R_1 (t > t_2) \quad (2.52)$$

3. Global attractor

3.1 the existence and uniqueness of solution

Theorem 3.1 Assume $(G_1), (G_2), (G_3)$, holds, and Lemma1 Lemma2 holds; the problem (1.1)-(1.3) exists a unique



smooth solution

$$(u, v) \in L^\infty([0, +\infty); H^{2m} \times H^m) \quad (3.1)$$

Proof. By the Galerkin method, Lemma 1 and Lemma 2, we can easily obtain the existence of solution. Next, we prove the uniqueness of Solutions in detail

Assume u, v are two solutions of the problems(1.1)-(1.3). let $w = u - v$, then

$$w(x, 0) = w_0(x) = 0, \quad w_t(x, 0) = w_1(x) = 0$$

Then two equations subtract and obtain

$$w_{tt} + (-\Delta)^m w_t + \left\| \nabla^m u \right\|^{2q} (-\Delta)^m u - \left\| \nabla^m v \right\|^{2q} (-\Delta)^m v + h(u_t) - h(v_t) = 0 \quad (3.2)$$

By multiplying (3.2) by w_t we get

$$(w_{tt} + (-\Delta)^m w_t + \left\| \nabla^m u \right\|^{2q} (-\Delta)^m u - \left\| \nabla^m v \right\|^{2q} (-\Delta)^m v + h(u_t) - h(v_t), w_t) = 0 \quad (3.3)$$

$$(w_{tt}, w_t) = \frac{1}{2} \frac{d}{dt} \|w_t\|^2 \quad (3.4)$$

$$(-\Delta)^m w_t, w_t) = \left\| \nabla^m w_t \right\|^2 \quad (3.5)$$

$$\begin{aligned} & (\left\| \nabla^m u \right\|^{2q} (-\Delta)^m u - \left\| \nabla^m v \right\|^{2q} (-\Delta)^m v, w_t) \\ &= (\left\| \nabla^m u \right\|^{2q} (-\Delta)^m u - \left\| \nabla^m u \right\|^{2q} (-\Delta)^m v + \left\| \nabla^m u \right\|^{2q} (-\Delta)^m v - \left\| \nabla^m v \right\|^{2q} (-\Delta)^m v, w_t) \\ &= (\left\| \nabla^m u \right\|^{2q} (-\Delta)^m w, w_t) + ((\left\| \nabla^m u \right\|^{2q} - \left\| \nabla^m v \right\|^{2q}) (-\Delta)^m v, w_t) \end{aligned} \quad (3.6)$$

$$\frac{1}{2} \frac{d}{dt} (\left\| \nabla^m u \right\|^{2q} \left\| \nabla^m w \right\|^2) - q \left\| \nabla^m u \right\|^{2q-1} \left\| \nabla^m u \right\| \left\| \nabla^m w \right\|^2$$

$$+ \left\| \nabla^m u \right\|^{2q} (-\Delta)^m v - \left\| \nabla^m v \right\|^{2q} (-\Delta)^m v, w_t)$$

$$\left| ((\left\| \nabla^m u \right\|^{2q} - \left\| \nabla^m v \right\|^{2q}) (-\Delta)^m v, w_t) \right|$$

$$\leq 2q \left\| \nabla^m u \right\| + \theta (\left\| \nabla^m v \right\| - \left\| \nabla^m u \right\|)^{2q-1} \left\| \nabla^m w \right\| \left\| (-\Delta)^m v \right\| \left\| w_t \right\|. \quad (3.7)$$

According to Lemma1, Lemma2,we have

$$2q \left\| \nabla^m u \right\| + \theta (\left\| \nabla^m v \right\| - \left\| \nabla^m u \right\|)^{2q-1} \left\| \nabla^m w \right\| \left\| (-\Delta)^m v \right\| \leq C_7,$$

(3.8)



$$q \left\| \nabla^m u \right\|^{2q-1} \left\| \nabla^m u_t \right\| \leq C_8.$$

Then

$$\left| ((\left\| \nabla^m u \right\|^{2q} - \left\| \nabla^m v \right\|^{2q}) (-\Delta)^m v, w_t) \right| \leq C_7 \left\| \nabla^m w \right\| \|w_t\|. \quad (3.9)$$

According to Young's inequality, we get

$$\left| ((\left\| \nabla^m u \right\|^{2q} - \left\| \nabla^m v \right\|^{2q}) (-\Delta)^m v, w_t) \right| \leq \frac{C_7}{2} \left\| \nabla^m w \right\|^2 + \frac{C_7}{2} \|w_t\|^2. \quad (3.10)$$

From (G_3) , we have

$$\left| (h(u_t) - h(v_t), w_t) \right| = \left| (h'(s) w_t, w_t) \right| \leq C_2 \|w_t\|^2. \quad (3.11)$$

From above, we have

$$\begin{aligned} & \frac{d}{dt} (\|u_t\|^2 + \left\| \nabla^m u \right\|^{2q} \left\| \nabla^m w \right\|^2) + 2 \left\| \nabla^m w_t \right\|^2, \\ & - (C_7 + 2C_8) \left\| \nabla^m w \right\|^2 - (C_7 + 2C_2) \|w_t\|^2 \leq 0. \end{aligned} \quad (3.12)$$

Then

$$\frac{d}{dt} (\|u_t\|^2 + \left\| \nabla^m u \right\|^{2q} \left\| \nabla^m w \right\|^2) \leq (C_7 + 2C_8) \left\| \nabla^m u \right\|^2 + (C_7 + 2C_2) \|w_t\|^2. \quad (3.13)$$

According to $\left\| \nabla^m u \right\|^{2q} \left\| \nabla^m w \right\| \geq \varepsilon \left\| \nabla^m w \right\|^2$, then

$$\frac{d}{dt} (\|u_t\|^2 + \left\| \nabla^m u \right\|^{2q} \left\| \nabla^m w \right\|^2) \leq \left(\frac{C_7 + 2C_8}{\varepsilon} \right) \left\| \nabla^m u \right\|^{2q} \left\| \nabla^m w \right\|^2 + (C_7 + 2C_2) \|w_t\|^2 \quad (3.14)$$

Taking $\gamma = \max\{\frac{C_7 + 2C_8}{\varepsilon}, C_7 + 2C_2\}$, we have

$$\frac{d}{dt} (\|w_t\|^2 + \left\| \nabla^m u \right\|^{2q} \left\| \nabla^m w \right\|^2) \leq \gamma (\left\| \nabla^m u \right\|^{2q} \left\| \nabla^m w \right\|^2 + \|w_t\|^2) \quad (3.15)$$

By using Gronwall inequality, we obtain

$$\|w_t\|^2 + \left\| \nabla^m u \right\|^{2q} \left\| \nabla^m w \right\|^2 \leq \gamma (\left\| \nabla^m u \right\|^{2q} \left\| \nabla^m w(0) \right\|^2 + \|w_t(0)\|^2) e^{\gamma t}. \quad (3.16)$$

Therefore

$$u = v$$

So we get the uniqueness of the solution.

3.2 Global attractor



Theorem 3.1. [11] Let E_1 be a Banach space, and $\{S(t)\}_{t \geq 0}$ are the semigroup operator on E_1 .

$S(t) : E_1 \rightarrow E_1$, $S(t+s) = S(t)S(s)$ ($\forall t, s \geq 0$), $S(0)=I$, where I is a unit operator. set $S(t)$ satisfy the follow conditions.

1) $S(t)$ is uniformly bounded, namely $\forall R > 0$, $\|u\|_{E_1} \leq R$, it exists a constant $C(R)$, so that

$$\|S(t)u\|_{E_1} \leq C(R) \quad (t \in [0, +\infty));$$

2) It exists a bounded absorbing set $B_0 \subset E_1$, namely, $\forall B \subset E_1$, it exists a constant t_0 , so that $S(t)B \subset B_0$ ($t \geq t_0$);

Where B_0 and B are bounded sets.

3) When $t > 0$, $S(t)$ is a completely continuous operator A .

Therefore, the semigroup operator $S(t)$ exists a compact global attractor.

Theorem 3.2. [12] Under the assume of Lemma 1, Lemma 2, Theorem 3.1, equations have global attractor

$$A = \omega(B_0) = \overline{\bigcap_{s \geq 0} \bigcup_{t \geq s} S(t)B_0}.$$

Where $B_0 = \{(u, v) \in H^{2m}(\Omega) \times H_0^m(\Omega) : \|u, v\|_{H^{2m} \times H_0^m}^2 = \|u\|_{H^{2m}}^2 + \|v\|_{H_0^m}^2 \leq R_0 + R_1\}$, B_0 is the bounded absorbing set

of $H^{2m}(\Omega) \times H_0^m(\Omega)$ and satisfies

$$(1) \quad S(t)A = A, \quad t > 0;$$

$$(2) \quad \lim_{t \rightarrow \infty} \text{dist}(S(t)B, A) = 0, \quad \text{here } B \subset H^{2m} \times H_0^m \text{ and it is a bounded set},$$

$$\lim_{t \rightarrow \infty} \text{dist}(S(t)B, A) = \sup_{x \in B} (\inf_{y \in A} \|S(t)x - y\|_{H^{2m} \times H_0^m}) \rightarrow 0, \quad t \rightarrow \infty.$$

Proof .Under the conditions of Theorem 3.1, it exists the solution semigroup $S(t)$,

$$S(t) : H^{2m}(\Omega) \times H_0^m(\Omega) \rightarrow H^{2m}(\Omega) \times H_0^m(\Omega), \quad \text{here } E_1 = H^{2m}(\Omega) \times H_0^m(\Omega).$$

(1) from Lemma 2.1 to Lemma 2.2, we can get that $\forall B \subset H^{2m} \times H_0^m$ is a bounded set that includes in the ball

$$\{\|u, v\|_{H^{2m} \times H_0^m} \leq R\},$$

$$\begin{aligned} \|S(t)(u_0, v_0)\|_{H^{2m} \times H_0^m}^2 &= \|u\|_{H^{2m}}^2 + \|v\|_{H_0^m}^2 \leq \|u_0\|_{H^{2m}}^2 + \|v_0\|_{H_0^m}^2 + C \leq R^2 + C \\ (t \geq 0, u_0, v_0 \in B) \end{aligned}$$

This shows that $S(t)$ ($t \geq 0$) is uniformly bounded $H^{2m}(\Omega) \times H_0^m(\Omega)$.

(2) Furthermore, for any $(u_0, v_0) \in H^{2m}(\Omega) \times H_0^m(\Omega)$, when $t \geq \max\{t_1, t_2\}$, we have,



$$\|S(t)(u_0, v_0)\|_{H^{\frac{2m}{2m+1}} \times H^{\frac{m}{2m+1}}}^2 = \|u\|_{H^{2m}(\Omega)}^2 + \|v\|_{H_0^m(\Omega)}^2 \leq R_0 + R_1$$

So we get B_0 is the bounded absorbing set.

(3) Since $E_1 = H^{2m}(\Omega) \times H^m(\Omega) \rightarrow E_0 = H^m(\Omega) \times L^2(\Omega)$ is compact embedded, which means that the bounded set in E_1

is the compact set in E_0 , so the semigroup operator $S(t)$ exists a compact global attractor A .

4 The estimates of the upper bounds of Hausdorff and fractal dimensions for the global attractor

We rewrite the problems (1.1)-(1.3):

$$u_{tt} + A^m u_t + \left\| A^{\frac{m}{2}} u \right\|^{2q} A^m u + h(u_t) = f(x). \quad (4.1)$$

Let $Au = -\Delta u$, where Ω is a bounded domain in R^N with smooth boundary $\partial\Omega$. The linearized equations of the above equation as follows

$$U_{tt} + AU = FU, \quad (4.2)$$

$$U_0 = \xi, U_t(0) = \zeta. \quad (4.3)$$

Let $U_0 \in H_0^m(\Omega)$, $U(t)$ is the solution of problems (4.20)-(4.21). We can prove that the problems (4.20)-(4.21) have a

unique solution $U \in L^\infty(0, T, H_0^m(\Omega))$, $U_t \in L^\infty(0, T, L^2(\Omega))$. The equation (4.20) is the linearized equation by the

equation (4.17). Define the mapping $Ls(t)_{u_0} : Ls(t)_{u_0} \zeta = U(t)$, here $u(t) = s(t)u_0$, let $\varphi_0 = (u_0, u_1)$,

$$\overline{\varphi_0} = \varphi_0 + \{\xi, \zeta\} = \{u_0 + \xi, u_1 + \zeta\}, \text{let :}$$

$$\|\varphi_0\|_{E_0} \leq R_1, \|\overline{\varphi_0}\|_{E_0} \leq R_2, E_0 = \nu \times H, E_0 = \nu \times H, \nu := H_0^m(\Omega), H := L^2(\Omega), S(t)\varphi_0 = \varphi(t) = \{u(t), u_t(t)\},$$

$$S(t)\overline{\varphi_0} = \{\varphi(t), \overline{\varphi_t}(t)\}$$

Lemma 4.1^[9] Assume H is a Hilbert space, E_0 is a compact set of H . $S(t) : E_0 \rightarrow H$ is a continuous mapping, satisfy the follow conditions.

$$1); S(t)E_0 = E_0, t > 0;$$

$$2) \text{ if } S(t) \text{ is Fréchet differentiable, it exists is a bounded linear differential operator } L(t, \varphi_0) \in C^+(R^+; L(E_0, E_0)),$$

$\forall t > 0$. that is



$$\frac{\left\| S(t)\overline{\varphi_0} - S(t)\varphi_0 - L(t,\varphi_0)(u,v) \right\|_{E_0}^2}{\left\| \{\xi, \zeta\} \right\|_{E_0}^2} \rightarrow 0, \{\xi, \zeta\} \rightarrow 0$$

The proof of lemma 4.1 see ref.[9],is omitted here .According to Lemma 4.1,we can get the following theorem:

Theorem 4.1^[9-10] Let A is the global attractor that we obtain in section 3.In that case , A has finite Hausdorff dimensions and fractal dimensions in $(H^{2m}(\Omega) \cap H_0^m(\Omega)) \times H_0^m(\Omega)$,that is $d_H(A) \leq n$ $d_F(A) \leq \frac{6n}{5}$.

Proof. Firstly, we rewrite the equations (4.17),(4.18) into the first order abstract evolution equations in E_0 ..

Let $\Psi = R_\varepsilon \varphi = \{u, u_t + \varepsilon u\}$,let $R_\varepsilon : \{u, u_t\} \rightarrow \{u, u_t + \varepsilon u\}$,is an isomorphic mapping.So let is the global attractor of $\{S(t)\}$,then $R_\varepsilon A$ is also the global attractor of $\{S_\varepsilon(t)\}$,then Ψ satisfies as follows:

$$\Psi_t + \Lambda_\varepsilon \Psi + \overline{h}(\Psi) = \overline{f}, \quad (4.4)$$

$$\Psi(0) = \{u_0, u_1 + \varepsilon u_0\}^T \quad (4.5)$$

Where $\Psi = \{u, u_t + \varepsilon u\}^T$, $\overline{h}(\Psi) = \{0, h(u_t)\}^T$, $\overline{f} = \{0, f(x)\}^T$

$$\Lambda_\varepsilon = \begin{pmatrix} \varepsilon I & -I \\ \left\| A^{\frac{m}{2}} u \right\|^{2q} & A^m - \varepsilon A^m + \varepsilon^2 I \\ A^m - \varepsilon I & A^m - \varepsilon I \end{pmatrix} \quad (4.6)$$

$$\Psi_t := F(\Psi) = \overline{f} - \Lambda_\varepsilon \Psi - \overline{h}(\Psi) \quad (4.7)$$

$$P_t = F_t(\Psi) \quad (4.8)$$

$$P_t + \Lambda_\varepsilon P + \overline{h}(\Psi) = 0 \quad (4.9)$$

Where $P = \{U, U_t + \varepsilon U\}^T$, $\overline{h}_t(\Psi)P = \{0, h_t(u_t)U\}^T$ The initial condition (4.3) can be written in the following form:

$$P(0) = w, w = \{\xi, \zeta\} \in E_0. \quad (4.10)$$

We take $n \in N$,then consider the corresponding n solution: $(P = P_1, P_2 \dots P_n, P_j \in E_0)$ of the initial values:

$(w = w_1, w_2 \dots w_n, w_j \in E_0)$ in the equations (4.8)-(4.10)

So there is $\left| P_1(t) \Lambda P_2(t) \Lambda \dots \Lambda P_n(t) \right|_{\Lambda_{E_0}^n} = \left| w_1 \Lambda w_2 \dots \Lambda w_n \right|_{\Lambda_{E_0}^n} \cdot e^{\int_0^t TrF_t(S_\varepsilon(\tau)\Psi_0)Q_n(\tau)d\tau}$



Form $\psi(\tau) = S_\varepsilon(\tau)\Psi_0$, we get $S_\varepsilon(\tau):\{u_0, v_1 = u_1 + \varepsilon u_0\} \rightarrow \{u(\tau), v(\tau) = u_t(\tau) + \varepsilon u(\tau)\}$,

$\psi(\tau) = \{u(\tau), v_t(\tau) = u_t(\tau) + \varepsilon u(\tau)\}$, here u is the solution of problems (4.1); Λ represents the outer product,

Tr represents the trace, $Q_n(\tau) = Q_n(\tau, \Psi_0; w_1, w_2 \dots w_n)$ is an orthogonal projection from the space $E_0 = V \times H$ to

the subspace spanned by $\{P_1(\tau), P_2(\tau), \dots, P_n(\tau)\}$. For a given time τ , let $\phi_j(\tau) = \{\xi_j(\tau), \zeta_j(\tau)\}$, $j = 1, 2, \dots, n$.

$\{\phi_j(\tau)\}_{j=1,2,\dots,n}$ is the standard orthogonal basis of the space $Q_n(\tau)_{E_0} = \text{span}[P_1(\tau), P_2(\tau), \dots, P_n(\tau)]$.

From the above, we have

$$\begin{aligned} Tr F_t(\Psi(\tau)) \cdot Q_n(\tau) &= \sum_{j=1}^n F_t(\Psi(\tau)) \cdot Q_n(\tau) \phi_j(\tau), \phi_j(\tau) \Big|_{E_0} \\ &= \sum_{j=1}^n F_t(\Psi(\tau)) \phi_j(\tau), \phi_j(\tau) \Big|_{E_0} \quad (4.11) \end{aligned}$$

Where $(\cdot, \cdot)_{E_0}$ is the inner product in E_0 . Then $(\{\xi, \zeta\}, \{\bar{\xi}, \bar{\zeta}\})_{E_0} = (\xi, \bar{\xi}) + (\zeta, \bar{\zeta})$;

$$(F_t(\Psi)\phi_j, \phi_j)_{E_0} = -(\Lambda_\varepsilon \phi_j, \phi_j)_{E_0} - (h_t(u)\xi_j, \xi_j);$$

$$(\Lambda_\varepsilon \phi_j, \phi_j)_{E_0} = \varepsilon \|\xi_j\|^2 + ((\left\| A^{\frac{m}{2}} u \right\|^{2q} A^m - \varepsilon A^m + \varepsilon^2 I) \zeta_j, \xi_j)$$

$$- (\xi_j, \zeta_j) + (A^m - \varepsilon I)(\zeta_j, \zeta_j) \quad (4.12)$$

$$\geq a (\|\xi_j\|^2 + \|\zeta_j\|^2).$$

$$\text{Where } a := \min\left\{ \frac{2\varepsilon + \varepsilon^2 - (\left\| A^{\frac{m}{2}} u \right\|^{2q} - \varepsilon) \lambda_1 - 1}{2}, \frac{-2\varepsilon + \varepsilon^2 + (2 - \left\| A^{\frac{m}{2}} u \right\|^{2q} - \varepsilon) \lambda_1 - 1}{2} \right\}.$$

Now suppose that $\{u_0, u_1\} \in A$, according to theorem 3.3, A is a bounded absorbing set in E_1 .

$$\Psi(t) = \{u(t), u_t(t) + \varepsilon u(t)\} \in D(A); D(A) = \{u \in V, Au \in H\}.$$

Then there is a $s \in [0, 1]$ to make the mapping $h_t : D(A) \rightarrow \rho(V_s, H)$. At the same time, there are the following results:



$$R_A = \sup_{\{\xi, \zeta\} \in A} |A \xi| < \infty$$

$$\sup_{\substack{u \in D(A) \\ |Au| < R_A}} |h_t(u)|_{\rho(v_s, H)} \leq r < \infty \quad (4.13)$$

Where $\|h_t(u)\xi_j, \zeta_j\|$ meets: $\|h_t(u)\xi_j, \zeta_j\| \leq r \|\xi_j\|_s \|\zeta_j\|$, Comprehensive above can be obtained

$$\begin{aligned} (F_t(\Psi)\phi_j, \phi_j)_{E_0} &\leq -a(\|\xi_j\|^2 + \|\zeta_j\|^2) + r \|\xi_j\|_s \|\zeta_j\| \\ &\leq -\frac{a}{2}(\|\xi_j\|^2 + \|\zeta_j\|^2) + \frac{r^2}{2a} \|\xi_j\|_s^2 \end{aligned} \quad (4.14)$$

Where $\|\xi_j\|^2 + \|\zeta_j\|^2 = \|\phi_j\|_{E_0}^2 = 1$, due to $\{\phi_j(\tau)\}_{j=1,2,\dots,n}$ is a standard orthogonal basis in $Q_n(\tau)_{E_0}$

So

$$\sum_{j=1}^n (F_t(\Psi(\tau))\phi_j(\tau), \phi_j(\tau))_{E_0} \leq -\frac{na}{2} + \frac{r^2}{2a} \|\xi_j\|_s^2 \quad (4.15)$$

Almost to all t , making

$$\sum_{j=1}^n \|\xi_j\|_s^2 \leq \sum_{j=1}^{n-1} \lambda_j^{s-1} \quad (4.16)$$

So

$$Tr F_t(\Psi(\tau)) \cdot Q_n(\tau) \leq -\frac{na}{2} + \frac{r^2}{2a} \sum_{j=1}^{n-1} \lambda_j^{s-1} \quad (4.17)$$

Let us assume that $\{u_0, u_t\} \in A$, is equivalent to $\Psi_0 = \{u_0, u_1 + \varepsilon u_0\} \in R_\varepsilon A$

Then

$$q_n(t) = \sup_{\Psi_0 \in R_\varepsilon A} \sup_{\substack{w \in E_0 \\ \|w\|_{E_0} \leq 1}} (\int_0^t Tr F_t(S_\varepsilon(\tau)\Psi_0) \cdot Q_n(\tau) d\tau) . j = 1, 2, \dots, n.$$

$$q_n = \lim_{t \rightarrow \infty} \sup q_n(t) \quad (4.18)$$

According to (4.17), (4.18), so

$$q_n(t) \leq -\frac{na}{2} + \frac{r^2}{2a} \sum_{j=1}^{n-1} \lambda_j^{s-1}$$



$$q_n \leq -\frac{na}{2} + \frac{r}{2a} \sum_{j=1}^{n-1} \lambda_j^{s-1} \quad (4.19)$$

Therefore, the Lyapunov exponent of A (or $R_\varepsilon A$) is uniformly bounded

$$\mu_1 + \mu_2 + \dots + \mu_n \leq -\frac{na}{2} + \frac{r^2}{2a} \sum_{j=1}^n \lambda_j^{s-1} \quad (4.20)$$

From what has been discussed above, it exists $n > 1$ and $s \in [0, 1]$, a, r are constants, then

$$\frac{1}{n} \sum_{j=1}^n \lambda_j^{s-1} \leq \frac{a^2}{6r^2} \quad (4.21)$$

$$q_n \leq -\frac{na}{2} \left(1 - \frac{r^2}{a^2} \sum_{j=1}^n \lambda_{j-1}^{s-1}\right) \leq -\frac{5na}{12} \quad (4.22)$$

$$(q_j)_+ \leq -\frac{r^2}{2a} \sum_{i=1}^j \lambda_i^{s-1} \leq \frac{na}{12}, \quad j = 1, 2, \dots, n. \quad (4.23)$$

$$\max_{1 \leq j \leq n-1} \frac{(q_j)_+}{|q_m|} \leq \frac{1}{5} \quad (4.24)$$

So, we immediately to the Hausdorff dimension and fractal dimension are respectively

$$d_H(A) < n, \quad d_F(A) < \frac{6}{5}n.$$

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6 References

- [1] Nakao M, Zhijian Y. D.Gobal attractors for some quasi-linear wave equations with a strong dissipation. *Advan Math Sci App*; 2007, 17:89-105.
- [2] Xiaoming F, Shengfan Z. Kernel sections for non-autonomous strongly damped wave equation of non-degenerate Kirchhoff-type. *Applied Mathematics and Computation*, 2004, 158:253-266.
- [3] Zhijian Yang, Baoxia Jin. Global attractor for a class of Kirchhoff models. *Mathematical Physics*. 2009.
- [4] Salim A, Messaoudi, Belkacem Said Houari. A blow-up result for a higher-order nonlinear Kirchhoff-type hyperbolic equation. *Appiled Mathematics letters*. 20(2007) 866-871.
- [5] Fucai Li Global Existence and Blow-up of Solutions for a Higher-order Kirchhoff-type Equation with Nonlinear Dissipation. *Applied Mathematics letters* 17(2004) 1409-1414.
- [6] Yang Zhijian and Wang Yunqing. Global attractor for the Kirchhoff type equation with a stromg dissipation. *Journal of Differential Equations*. 249(2010), pp.3258-3278.
- [7] ZhijianYang, Jianling Cheng. Asymptotic Behavior of Solution to the Kirchhoff Type Equation. *Mathematical Physics*. 2011, 31A(4):1008-1021 .



- [8] Meixia Wang; Cuicui Tian; Guoguang Lin. Global attractor and dimension estimation for a 2D generalized anisotropy Kuramoto-Sivashinsky equation. International Journal of Modern Nonlinear Theory and Application.3(2014),pp/163-172.
- [9] Teman,R. Infinite Dimensional Dynamics Systems in Mechanics and Physics.Springer, New York(1998).
- [10] Jingzhuwu, Guoguang Lin. The global attractor of the Boussinesq equation with damping term and its dimension estimation.Journal of Yunnan University, 31(2009),pp.335-340
- [11] Lin Guoguang, Nonlinear evolution equation, Yunnan University Press, 2011.